

Density-Functional-Theory Lecture I

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Overview: Monday

- Building the periodic table:
 - Classical mechanics.
 - Electricity.
 - Quantum mechanics: wave function and probability.
 - Hydrogen atom/units.
 - Multi-electron atoms.
 - Periodic table.
- Covalent bonding.
- Adiabatic-decoupling/time scales.
- Ritz variational principle.
- Hartree-Fock: Exchange Interaction.
- Density-Functional-Theory.
- Exchange-correlation energy.
- Examples of DFT applications.

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Material Properties

- Crystal Chemistry.
- Ground State Energy.
- Phase Diagrams.
- Vibrational Properties.
- Electronic Density of States.
- Magnetism.
- Solid Solutions. Alloying.
- Elasticity.
- ...



- Transport Properties:
- Thermal conductivity.
 - Electrical conductivity.
 - Viscosity.
 - ...

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"Minimal" Requirements

- Periodic table.
- Thermodynamics.
- Single theoretical framework for solid, liquid, gas.
- Bulk, surfaces, and molecules.



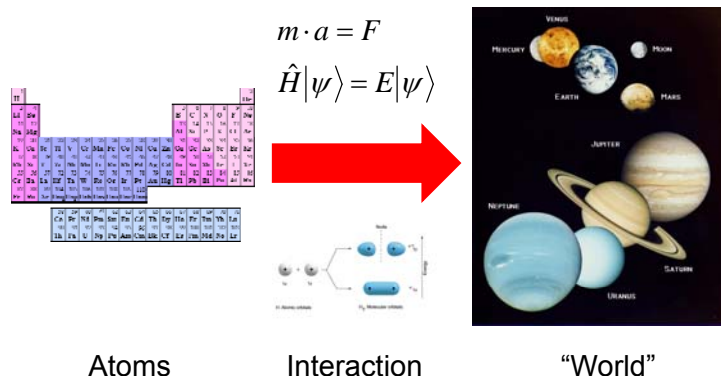
Requirement

- Description of atoms and interactions.
- Forces.
- Ensembles.

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Scientific Goal



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Classical Mechanics, 18th Century



Sir Isaac Newton
(1643 – 1727)



$$m \cdot \ddot{\vec{x}}_i = m \cdot \vec{a}_i = \vec{F}_i$$

$$\vec{F}_i = -\vec{\nabla}_i V(\{\vec{r}_j\})$$

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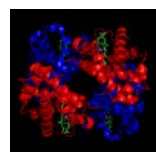
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Empirical Potentials

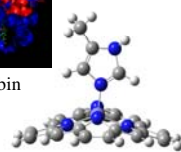
Born-Meyer potential:

$$V(r) = \frac{Z_1 Z_2}{r} + A \exp(-B \cdot r) + \frac{C}{r^6}$$

Coulomb-Potential Non-Coulomb Dispersion



Hemoglobin



Active site: Fe²⁺

Newton's 2nd law:

$$m \cdot \vec{a}_i = \vec{F}_i = \frac{1}{2} \sum_{i \neq j} \nabla_i V(\vec{r}_j)$$

Advantage:

- Fast.
- Many atoms (> Millions).
- Explicit dispersion forces.

Disadvantage:

- Charges (Z); Constants: A, B, C.
- Multiple valance states: 3d, 4d, ... Fe²⁺, Fe³⁺, ...

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http://upload.wikimedia.org/wikipedia/commons/3/3d/1GZX_Haemoglobin.png

Assessment

- Transferability?
 - Electronic properties?
 - Predictive Power?
- Suggest that electrons should be treated explicitly

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Electricity, 18th+19th Century



Charles Coulomb
(1736 – 1806)

$$F = -\frac{1}{4\pi\epsilon_0} \frac{q_1 \cdot q_2}{R^2}$$



James Clerk Maxwell
(1831 – 1879)

Electrodynamics

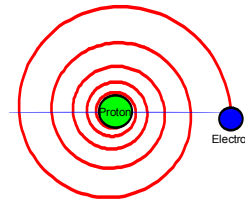
$$\begin{aligned}\nabla \times \vec{E} &= -\frac{\partial \vec{B}}{\partial t} - \vec{M} \\ \nabla \times \vec{H} &= -\frac{\partial \vec{D}}{\partial t} + \vec{J} \\ \nabla \cdot \vec{D} &= \rho \\ \nabla \cdot \vec{B} &= 0\end{aligned}$$

<http://www.microwaves101.com/encyclopedia/images/maxwellsequations/maxwellsequations.jpg>
<http://upload.wikimedia.org/wikipedia/commons/0/08/08+08/09/2009, UNM>
http://outreach.jach.hawaii.edu/presroom/2008_jcm/day/Maxwell.jpg

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Constituents of Atoms

- Electron (Thomson, 1897):
negatively charged particle.
→ need positively charged particle.
- Proton (Thomson + Rutherford, 1907).

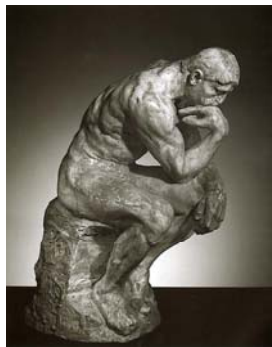


Electrodynamics predicts that the electron spirals into the nucleus within $\sim 10^{-8}$ s.

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Classical Physics Predicts that Matter is Unstable?!?



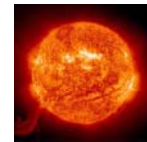
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Auguste Rodin: The Thinker

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Solution: Quantum Mechanics

➤ Max Planck (1899).



$$E = n\hbar\omega$$

➤ Albert Einstein (1905).

➤ Niels Bohr (1915).

➤ Werner Heisenberg (1925).

➤ Erwin Schroedinger (1925).

$$i\hbar \frac{\partial \psi}{\partial t} = \hat{H}\psi$$

$$E\psi = \hat{H}\psi$$

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Interpretation

(Time independent) Schrodinger equation:

Energy eigenvalue Total energy (operator)

$$E\psi = \hat{H}\psi$$

New quantity: Wave function

$$\hat{H} = \hat{K} + \hat{U}$$

Total energy = Kinetic energy + Potential energy

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What is a Wave Function?

Interpretation (Max Born, 1927):

probability

$$= |\psi|^2 dV$$

$$= \langle \psi | \psi \rangle dV$$

The square of a wave function describes a probability (density).

Example: Hydrogen Atom

$$\hat{K} = -\frac{\hbar^2}{2m} \nabla^2; \hat{U} = -\frac{e^2}{r}$$

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Hydrogen Atom

$$E_n = -\frac{13.6\text{eV}}{n^2}$$

$$a_n = 0.529 \cdot 10^{-10} m$$

$$\hat{H} = \frac{\hat{p}^2}{2\mu} - \frac{e^2}{|\hat{r}|}$$

$$l = 0, 1, \dots, n-1$$

$$m_l = -l, \dots, l$$

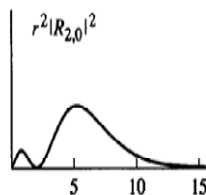
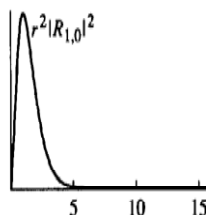
Degeneracy:

No-spin: n^2

Spin: $2 \cdot n^2$

$$1 \text{ eV} = 1.6022 \cdot 10^{-19} \text{ J} \\ = 96.4860 \text{ kJ/mol.}$$

$$1 \text{ kcal/mol} = 4.184 \text{ kJ/mol}$$



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Thermodynamics

How to distribute two or more electrons over available states?

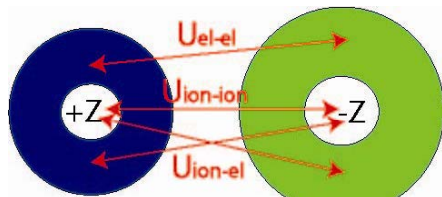
Lowest energy that is compatible with constraints

==
Ground State

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Interactions



Kinetic energies:
 K_{ion}
 K_{el}

Thermodynamics

Finite T : Minimize $F=F(V,T,N)$ or $G=G(P,T,N)$...
 T=0 K : Minimize $E=E(V,N)$ or $H=H(P,N)$

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Multi-Electron Atoms

Electrons are indistinguishable



Wolfgang Pauli:
 Pauli exclusion
 Principle (1925)



Physics/Chemistry
 must be independent
 of numbering.

In quantum mechanics we have two possibilities:

$$\psi(r_1, r_2) = -\psi(r_2, r_1)$$

Fermions: electrons,
 protons, neutrons,...

$$\psi(r_1, r_2) = +\psi(r_2, r_1)$$

Bosons: photons, phonons,
 spin-waves,...

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<http://www.csu.edu/~mdutton/u61s2s.jpeg>

Pauli-Exclusion-Principle

$$\text{Probability} = |\psi(r_1, r_2)|^2 dV_1 dV_2$$

Two Fermions in the same location:

$$\psi(r_1, r_2) = -\psi(r_2, r_1)$$

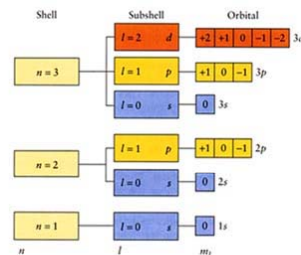
$$\Rightarrow \psi(r_1, r_1) = 0$$

$$\Rightarrow \text{Probability} = 0$$

Slater determinants:

$$\psi(r_1, \dots, r_n) = \frac{1}{N!} \begin{vmatrix} \psi_1(r_1) & \dots & \psi_1(r_n) \\ \dots & \dots & \dots \\ \psi_n(r_1) & \dots & \psi_n(r_n) \end{vmatrix}$$

Fermions in the exact same quantum state are impossible



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Hund's Rule #1

For a given electron configuration, the term with maximum multiplicity has the lowest energy. Since multiplicity is equal to $2S+1$, this is also the term with maximum S .

$$|1,2\rangle = \frac{1}{\sqrt{2}} (|\uparrow\rangle|\downarrow\rangle \mp |\downarrow\rangle|\uparrow\rangle) \quad * \quad \frac{1}{\sqrt{2}} (|1\rangle|2\rangle \pm |2\rangle|1\rangle)$$

"+" == symmetric; "-" == anti-symmetric.



Friedrich Hund

Consider two electrons:

$S=0 \rightarrow$ anti-symmetric spin state \rightarrow symmetric spatial state.

$S=1 \rightarrow$ symmetric spin state \rightarrow anti-symmetric spatial state
 \rightarrow zero probability to find electrons at "same" location.
 \rightarrow reduced Coulomb repulsion.

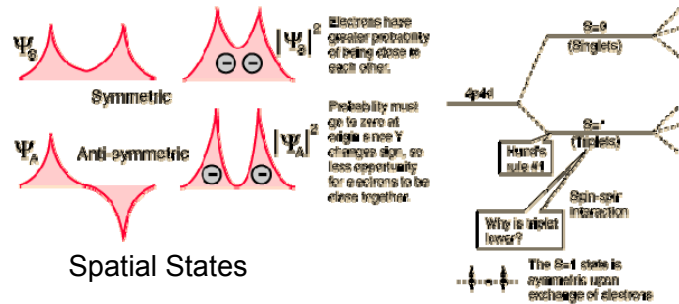
$\rightarrow S=1$ (triplet) state has a lower energy.

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<http://hyperphysics.phy-astr.gsu.edu/Hbase/Atomic/Hund.html#c2>

Hund's Rule #1, cont'd



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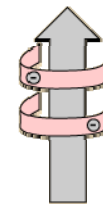
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Hund's Rule #2

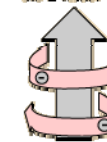
For a given multiplicity, the term with the largest value of L has the lowest energy.

High $L \rightarrow$ electrons orbit in the same direction. For low L some electrons rotate oppositely \rightarrow must pass \rightarrow increase in Coulomb repulsion \rightarrow higher energy.

High L , electrons orbiting same direction to add to L value.



Low L , some electrons orbiting in opposite direction to reduce the L value.



"Classically" speaking:
Larger L
 \rightarrow More electrons orbit nucleus in the same sense.
 \rightarrow Reduced Coulomb repulsion.
 \rightarrow Lower energy.

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<http://hyperphysics.phy-astr.gsu.edu/hbase/AtomicHund.html#c2>

Hund's Rule #3

For a given term, in an atom with outermost subshell half-filled or less, the level with the lowest value of

$$J=L+S$$

lies lowest in energy. If the outermost shell is more than half-filled, the level with highest value of J is lowest in energy.

Reason: Spin-orbit coupling $\sim L \cdot S \rightarrow J=L+S \rightarrow$ energy lower if L and S are in opposite directions \rightarrow lower J .

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Periodic Table

I	II	III	IV	V	VI	VII	VIII	IX	X	XI	XII	XIII	XIV	XV	XVI	XVII	XVIII
H 1 1s ¹																	
Li 3 2s ¹	Be 4 2s ²																
Na 11 3s ¹	Mg 12 3s ²																
K 19 4s ¹	Ca 20 4s ²	Sc 21 3d ¹ 4s ²	Ti 22 3d ² 4s ²	V 23 3d ³ 4s ²	Cr 24 3d ⁵ 4s ¹	Mn 25 3d ⁵ 4s ²	Fe 26 3d ⁶ 4s ²	Co 27 3d ⁷ 4s ²	Ni 28 3d ⁸ 4s ²	Cu 29 3d ¹⁰ 4s ¹	Zn 30 3d ¹⁰ 4s ²	Ga 31 4s ² 4p ¹	Ge 32 4s ² 4p ²	As 33 4s ² 4p ³	Se 34 4s ² 4p ⁴	Br 35 4s ² 4p ⁵	Kr 36 4s ² 4p ⁶
Rb 37 5s ¹	Sr 38 5s ²	Y 39 4d ¹ 5s ²	Zr 40 4d ² 5s ²	Nb 41 4d ⁴ 5s ¹	Mo 42 4d ⁵ 5s ¹	Tc 43 4d ⁵ 5s ²	Ru 44 4d ⁷ 5s ¹	Rh 45 4d ⁸ 5s ¹	Pd 46 4d ¹⁰ 5s ⁰	Ag 47 4d ¹⁰ 5s ¹	Cd 48 4d ¹⁰ 5s ²	In 49 5s ² 5p ¹	Sn 50 5s ² 5p ²	Sb 51 5s ² 5p ³	Te 52 5s ² 5p ⁴	I 53 5s ² 5p ⁵	Xe 54 5s ² 5p ⁶
Cs 55 6s ¹	Ba 56 6s ²	La 57 5d ¹ 6s ²	Ce 58 5d ¹ 6s ²	Pr 59 5d ¹ 6s ²	Nd 60 5d ¹ 6s ²	Pm 61 5d ⁴ 6s ²	Sm 62 5d ⁰ 6s ²	Eu 63 5d ⁰ 6s ²	Gd 64 5d ¹ 6s ²	Tb 65 5d ¹ 6s ²	Dy 66 5d ¹ 6s ²	Ho 67 5d ¹ 6s ²	Er 68 5d ¹ 6s ²	Tm 69 5d ¹ 6s ²	Yb 70 5d ¹ 6s ²	Lu 71 5d ¹ 6s ²	
Fr 87 7s ¹	Ra 88 7s ²	Ac 89 6d ¹ 7s ²	Th 90 6d ² 7s ²	Pa 91 5f ² 7s ²	U 92 5f ³ 7s ²	Np 93 5f ⁴ 7s ²	Pu 94 5f ⁶ 7s ²	Am 95 5f ⁷ 7s ²	Cm 96 5f ⁷ 7s ²	Bk 97 5f ⁷ 7s ²	Cf 98 5f ¹⁰ 7s ²	Es 99 5f ¹⁴ 7s ²	Fm 100 5f ¹⁴ 7s ²	Md 101 5f ¹⁴ 7s ²	No 102 5f ¹⁴ 7s ²	Lr 103 5f ¹⁴ 7s ²	
Lanthanides (Rare earths)																	
Actinides																	

$$N \rightarrow l-1 \rightarrow l_z = -l \dots +l = 2 * l + 1; s = \pm 1/2$$

$$N=1: l=0, s = \pm 1/2 \rightarrow 2 \text{ states}$$

$$N=2: l=0, s = \pm 1/2 \rightarrow 2 \text{ states}$$

$$l=1, s = \pm 1/2 \rightarrow 6 \text{ states}$$

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Summary

- Failure of classical mechanics/electrodynamics.
- Quantum mechanics to the rescue.
- New quantity: wave function
Probability.
- Building blocks of the periodic table:
Pauli exclusion principle.
Hund's rules.

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Born-Oppenheimer or Adiabatic Decoupling

Do we need to treat nuclei quantum mechanically as well?

$$\hat{H} = \hat{K}_{ion} + \hat{U}_{ion} + \hat{K}_{el} + \hat{U}_{el} + \hat{U}_{el-ion}$$

$m_p/m_e \sim 2000 \rightarrow$ motion of electron much faster than nuclei.

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Born-Oppenheimer or Adiabatic Decoupling

Time scales:

Diffusion : $D = 10^{-11} \text{ cm}^2/\text{s}$; $d = 1 \text{ \AA} \rightarrow t = d^2/D \sim 10^{-5} \text{ s}$

Vibrations : $H_2 \sim 4400 \text{ cm}^{-1}$ $t = 1/(c k) \sim 7.5 \cdot 10^{-15} \text{ s}$.

Electrons : $t = h/c^2 m_e \sim 10^{-20} \text{ s}$.

Electronic time-scale shortest

\rightarrow Electrons follow nuclear motion *instantaneously*.
(except may be for the lightest element: hydrogen)

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Consequence

If $T=0$ K \rightarrow Kinetic energy of the nuclei is zero.

Note: Nuclei are **not** eliminated from the problem, $U_{\text{nuc-nuc}}$ and $U_{\text{nuc-el}}$.
Static problem.

If $T>0$ K \rightarrow Nuclei are no longer fixed. Lattice vibrations; thermodynamics.

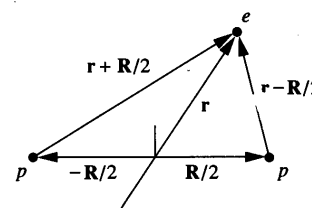
Electrons: always treated quantum mechanically.

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Covalent Bonds- I

$$\hat{H} = \frac{\hat{p}^2}{2m_e} - \frac{e^2}{|\hat{r} - R/2|} - \frac{e^2}{|\hat{r} + R/2|} + \frac{e^2}{|R|}$$



$$\hat{H}|1\rangle = \lambda|2\rangle \Rightarrow \lambda = \pm 1$$

$$|\pm\rangle = \frac{1}{\sqrt{2 \pm 2 \cdot \langle 1|2\rangle}} (|1\rangle \pm |2\rangle)$$

$$\hat{H}_{ij} = \langle i | \hat{H} | j \rangle$$

$$\hat{H}_{11} = \hat{H}_{22}$$

$$\hat{H}_{12} = \hat{H}_{21}$$

Electron bound to nucleus 1: $|1\rangle$
Electron bound to nucleus 2: $|2\rangle$

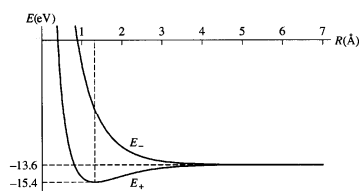
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Covalent Bonds- II

$$|\pm\rangle = \frac{1}{\sqrt{2 \pm 2 \cdot \langle 1|2\rangle}} (|1\rangle \pm |2\rangle)$$

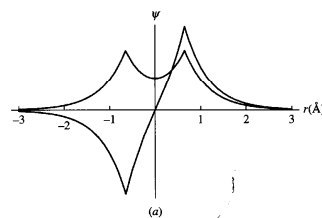
$$E_{\pm} = \frac{1}{1 \pm \langle 1|2\rangle} (H_{11} \pm H_{12})$$



$$\langle 1|2\rangle \neq 0$$

$$\Rightarrow H_{12} = \int \psi_1(r) H \psi_2(r) dV \neq 0$$

Overlap of electronic wave functions is necessary for covalent bonding.



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Ritz's Variational Principle

$$E_0 \leq \frac{\langle \Psi^*(r_1, \dots, r_N) | \hat{H} | \Psi(r_1, \dots, r_N) \rangle}{\langle \Psi^*(r_1, \dots, r_N) | \Psi(r_1, \dots, r_N) \rangle}$$

For a given H the ground state energy can be estimated as the lowest energy found for **any** possible wave function.

Also reminiscent of Pauli exclusion principle, Hund's rules, and thermodynamics which all required the minimization of energy.

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Generalization: Hartee-Fock

Many-electron systems:

Keep:

Adiabatic decoupling.

Pauli-principle.

$$\psi(r_1, \dots, r_n) = \frac{1}{N!} \begin{vmatrix} \psi_1(r_1) & \dots & \psi_1(r_n) \\ \dots & \dots & \dots \\ \psi_n(r_1) & \dots & \psi_n(r_n) \end{vmatrix}$$

This theory is called Hartree-Fock theory.

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Hartee-Fock: Electrostatics

Notably the theory contains expressions such as:

$$\iint \psi_i^2(r_1) \frac{1}{r_{12}} \psi_j^2(r_2) dV_1 dV_2$$

Electrostatic energy

Note:

$$\psi_i = \psi_j$$

Electrons at the same location: $r_{12}=0$
 \rightarrow Increase of electrostatic energy.

This is impossible according to the Pauli-exclusion principle.

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Hartee-Fock: Exchange

But another term appears as well:

$$-\iint \psi_i(r_1) \psi_j(r_2) \frac{1}{r_{12}} \psi_i(r_2) \psi_j(r_1) dV_1 dV_2$$

Exchange energy:

Origin: Pauli-exclusion principle.

No classical analogue.

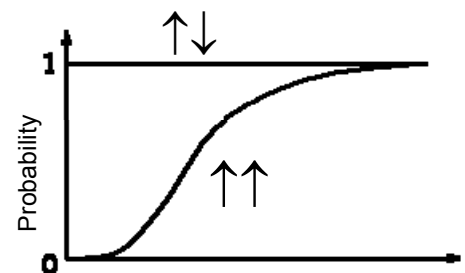
$$\psi_i = \psi_j$$

Exact compensation of electrostatic term:
 Pauli exclusion principle is fulfilled.

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Exchange: Another Look



Pauli exclusion principle

\rightarrow Exchange energy

\rightarrow Reduces Coulomb repulsion

\rightarrow Lowers energy of system.

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Changing the Point of view

Hartree-Fock:

central quantity: wave functions.

Alternatively: Choose to focus on the charge density.

→ Density-Functional-theory.

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Density-Functional-Theory (DFT)

Hohenberg-Kohn (1964)

Theorem 1

For any system of *interacting* particles in an external potential $V_{\text{ext}}(r)$, the potential is determined *uniquely* by the ground state particle density $n_0(r)$:

$$E[n, V_{\text{ext}}] = \int n(\vec{r}) \cdot V(\vec{r}) d^3r + F[n]$$

- $F[n]$ is unique and depends only on the density, n .
- Completely general.

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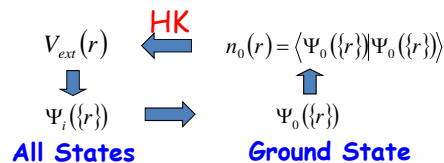
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Theorem 2

A *universal* functional for the energy $E[n, V_{\text{ext}}]$ in terms of the density $n(r)$ exists with a *global minimum* for the exact ground state density $n_0(r)$.

→ Valid for any system (universal): gas, liquid, solid...

Iterative Solution



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Kohn and Sham Ansatz (1965)

Replace *correlated* many-electron problem by an equivalent single electron *independent* particle problem in an *effective* potential.

$$F[n] = \frac{e^2}{2} \int \frac{n(\vec{r}) \cdot n(\vec{r}')}{|\vec{r} - \vec{r}'|} d^3r d^3r' \quad \text{Coulomb interaction}$$

$$+ T_0[n] \quad \text{Kinetic energy}$$

$$+ E_{\text{XC}}[n] \quad \text{Exchange-Correlation-energy}$$

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Kohn and Sham Equations

Kohn and Sham (1965)

$$\left[-\frac{\hbar^2}{2m_{el}} \nabla_r^2 + V_{eff}(r, n(r)) \right] \phi_i(r) = \varepsilon_i \phi_i(r)$$

$$V_{eff}(r, n(r)) = V(r)$$

External potential.

$$+ e^2 \int \frac{n(r')}{|r-r'|} dr'$$

Coulomb interaction.

$$+ \frac{\delta E_{xc}(n(r))}{\delta n(r)}$$

Exchange-correlation.

Can be generalized to include magnetism:

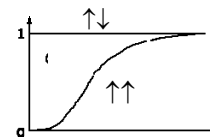
$$n = n_{\uparrow} + n_{\downarrow}$$

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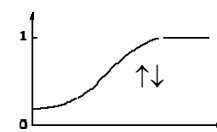
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What is the Exchange-Correlation Potential, E_{xc} ?

Thermodynamics: ground state == lowest energy state



Pauli exclusion principle
→ Exchange energy.



Repulsion of opposite spins
→ Correlation energy.

Both effects lower the energy

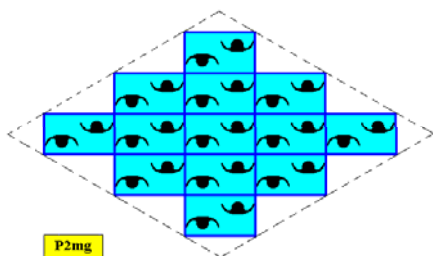
Active research field to develop better E_{xc} functionals:

- LDA : Ceperley and Alder.
- GGA : PW91, PBE, rPBE, revPBE,...

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Periodic Structures: Bloch's Theorem



Crystals are 3-d periodic structures:
That is they can be created by periodically repeating a smaller unit.

Crystal == Lattice + Motif

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http://www.metalphysics.nl/ontology/para_10.gif

Consequence?

1 g of Na → $\sim 3 \times 10^{23}$ electrons.

Need to solve determine for 10^{23} electrons?!?

Recognize that crystals are periodic structures.

Sufficient to describe one unit cell. (Bloch's theorem).

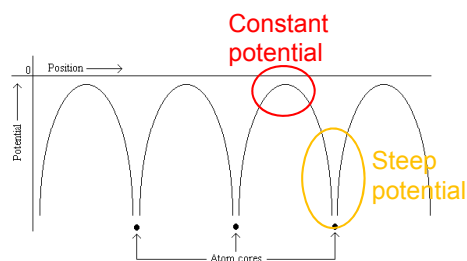
bcc-Na: 22 electrons sufficient.

→ Our problem has become tractable.

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Suitable Electronic Wave Functions



Solutions to Schrödinger equation with **constant potential**:

Planewaves: $\psi \sim \exp(i k r)$

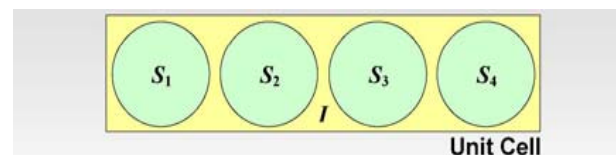
<http://gauss.admsmt.edu/~bhemmel/quantu1.gif>

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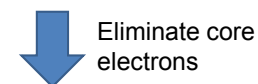
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Pseudopotentials

“Problem”: Steep potential close to nuclei
→ plane waves are not well suited.



Idea: Most material/chemical properties rely on valence electrons.



Pseudopotential.

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Optimization: Forces

Hellman-Feynman (force) theorem:

$$\begin{aligned}\vec{F}_i &= -\nabla_i E \\ &= \left\langle \psi \left| \nabla_i \hat{H}_{\text{electronic}} \right| \psi \right\rangle - \nabla_i U_{\text{ion-ion}}\end{aligned}$$

Summary:

- ✓ Periodic Table.
- ✓ Pseudopotential → fewer electrons.
- ✓ Universal E_{XC} : solid, liquid, gas. Bulk, surface, molecule.
- ✓ Optimization.

→ **Predictive power.**

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Summary

- Covalent bonding.
- Adiabatic decoupling.
T=0 K and finite temperatures.
- Hartree-Fock → Exchange.
- DFT → Exchange and correlation.
- Periodic structures, Bloch's theorem.
- Pseudopotential.
- Forces.
- Predictive power.

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Some Applications of DFT

“Fundamentals” of Condensed Matter Physics/Chemistry:

- Sodium at high pressure.
- Octet-rule at high pressure.

Student presentations

- Andrew: $\text{Pt}_{1-x}\text{Re}_x$ solid solutions.
- Eric: Pd on $\gamma\text{-Al}_2\text{O}_3$.
- Levi: Pd on $\alpha\text{-Al}_2\text{O}_3$.
- Sam: non-Pt based catalysts.

Sodium at High Pressures

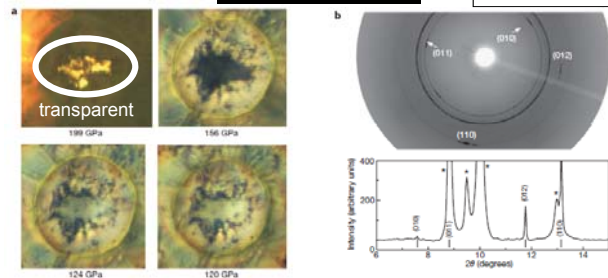
(Ma et al., 2009)

Background:

- Sodium at ambient conditions is a metal.
- Expected: structures adopt dense packed structures at sufficiently high pressure.

Experiment

1 GPa $\sim 10^4$ atm



Theory

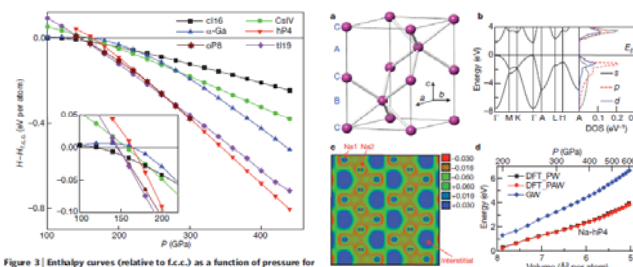


Figure 3 | Enthalpy curves (relative to f.c.c.) as a function of pressure for c16, c16V, α -Ga, α -Pb, t119 and hP4 structures of Na. The inset shows the transition sequence f.c.c. \rightarrow c16 \rightarrow t119 (α -Pb) \rightarrow hP4 in more detail. Two phase transitions, c16 \rightarrow t119 \rightarrow hP4, are predicted at 152 and 260 GPa, respectively. The α -Pb structure is less stable than the t119 phase in the pressure range 152–260 GPa, but is energetically very competitive. It is

- Phases: bcc \rightarrow fcc (close-packed) \rightarrow non-closed-packed...
- Transparent (non-metallic) high pressure phase.

Octet rule at High Pressures

(Kiefer and Tschauner)

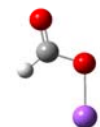
- Octet rule: main group elements tend to achieve a complete outermost shell: eight electrons.



• CH_4 : C: $4e^-$; each hydrogen donates $1e^-$.



• H_2S : S: $6e^-$; each hydrogen donates $1e^-$.



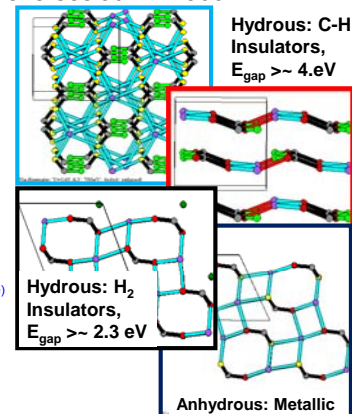
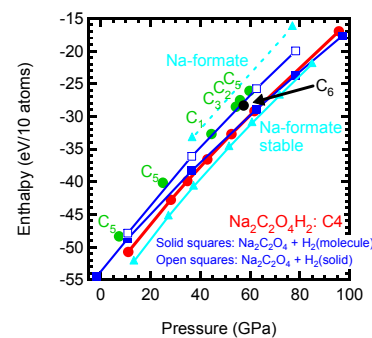
• CHOO-Na : Na-formate.

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Experimental Observations

- Sample becomes opaque at ~ 20 GPa during laser heating to $T \sim 1500$ K, reverses at $T \sim 1300$ K.
- H_2 libron observed.



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Summary

- DFT versatile tool.
- Exchange and correlation.
- Pseudopotentials.
- Periodic table.
- Applicable to solid, liquid, and gas.
- Bulk, surface, molecule.
- $T=0$ K, finite temperature.

Outlook

- Review.
- DFT: in's and out's.

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